

Determination of the b -quark leading shape function parameters in the shape function scheme using the Belle $B \rightarrow X_s \gamma$ photon energy spectrum

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Abstract

We determine the b -quark shape function parameters in the shape function scheme, $m_b(\text{SF})$ and $\mu_\pi^2(\text{SF})$ using the Belle $B \rightarrow X_s \gamma$ photon energy spectrum. We assume three models for the form of the shape function; exponential, gaussian and hyperbolic.

INTRODUCTION

One of the best ways to determine the Cabibbo-Kobayashi-Maskawa(CKM) matrix element $|V_{ub}|$ is to measure the inclusive charmless semileptonic ($B \rightarrow X_u \ell \nu$) partial branching fraction of the B meson. Its theoretical prediction relies on an accurate description of non-perturbative bound-state effects of a b quark in the B meson. The effects are encoded in B -meson shape functions, which are not theoretically calculable and have to be determined experimentally. The leading order shape function can be determined using the photon energy spectrum in $B \rightarrow X_s \gamma$ decays, since up to the leading order of $1/m_b$ in the Heavy Quark Expansion it is described with the same leading order shape function as the one used for the prediction of the $B \rightarrow X_u \ell \nu$ decays. Such shape function determination was first performed by CLEO on their $B \rightarrow X_s \gamma$ data [1], and was then repeated in Ref. [2] for the Belle $B \rightarrow X_s \gamma$ data, where the theory proposed by Kagan and Neubert [3] was used to describe the photon energy spectrum.

Recently, Bosch, Lange, Neubert and Paz presented theoretical expressions for the triple differential $B \rightarrow X_u \ell \nu$ decay rates and for the $B \rightarrow X_s \gamma$ photon spectrum, which incorporate all known contributions and smoothly interpolate between the "shape-function region" of large hadronic energy and small invariant mass, and the "OPE region" (in which all hadronic kinematical variables scale with M_B) [4, 5]. The differential rate is given in terms of a leading shape function and a subleading shape function in the shape function scheme [4, 5], where in this scheme the leading shape function is expressed with the parameters $m_b(\text{SF})$ and $\mu_\pi^2(\text{SF})$.

In this paper we report on the determination of the shape function parameters by fitting the Belle $B \rightarrow X_s \gamma$ data with the predicted photon energy spectrum [5], where the default subleading shape function model from Ref. [5] is used. It allows for a precise determination of $|V_{ub}|$ as shown in Ref. [6] and Ref. [7].

Apart from using an updated and more complete theoretical description of the $B \rightarrow X_s \gamma$ decays, this analysis differs from the one in Ref. [2] by using a different set of shape function model parametrizations, namely exponential, gaussian and hyperbolic. It also includes subleading shape function effects that were absent in the former analysis.

PROCEDURE

We used a method based on that devised by the CLEO Collaboration [8]. We fit Monte Carlo (MC) simulated spectra to the raw data photon energy spectrum. "Raw" refers to the spectra that are obtained after the application of the $B \rightarrow X_s \gamma$ analysis cuts. The use of "raw" spectra correctly accounts for Lorentz boost from the B rest frame to the center of mass system, energy resolution effects and avoids unfolding. The method is as follows:

1. Assume a shape function model.
2. Simulate the photon energy spectrum for a certain set of parameters; ($m_b(\text{SF})$, $\mu_\pi^2(\text{SF})$).
3. Perform a χ^2 fit of the simulated spectrum to the data where only the normalization of the simulated spectrum is floated and keep the resultant χ^2 value.
4. Repeat steps 2-3 for different sets of parameters to construct a two dimensional grid with each point having a χ^2 .

5. Find the minimum χ^2 on the grid and all points on the grid that are one unit of χ^2 above the minimum.
6. Repeat steps 1-5 for a different shape function model.

Shape function models

Three shape function forms suggested in the literature are employed; exponential, gaussian and hyperbolic [5]. Their functional forms are described in Table I: they are a function of $\hat{\omega}$ and are parameterized by two parameters: Λ and b . Example shape function forms are plotted in Fig. 1.

Shape Function	Form
exponential	$F^{(\text{exp})}(\hat{\omega}; \Lambda, b) = \frac{N^{(\text{exp})}}{\Lambda} \left(\frac{\hat{\omega}}{\Lambda}\right)^{b-1} \exp\left(-d_{\text{exp}} \frac{\hat{\omega}}{\Lambda}\right)$
gaussian	$F^{(\text{gauss})}(\hat{\omega}; \Lambda, b) = \frac{N^{(\text{gauss})}}{\Lambda} \left(\frac{\hat{\omega}}{\Lambda}\right)^{b-1} \exp\left(-d_{\text{gauss}} \frac{\hat{\omega}^2}{\Lambda^2}\right)$
hyperbolic	$F^{(\text{hyp})}(\hat{\omega}; \Lambda, b) = \frac{N^{(\text{hyp})}}{\Lambda} \left(\frac{\hat{\omega}}{\Lambda}\right)^{b-1} \cosh^{-1}\left(d_{\text{hyp}} \frac{\hat{\omega}}{\Lambda}\right)$

where the constants are:

$$\begin{aligned}
 N^{(\text{exp})} &= \frac{d_{(\text{exp})}^b}{\Gamma(b)}, & d_{(\text{exp})} &= b, \\
 N^{(\text{gauss})} &= \frac{2 d_{(\text{gauss})}^{b/2}}{\Gamma(b/2)}, & d_{(\text{gauss})} &= \left(\frac{\Gamma(\frac{1+b}{2})}{\Gamma(\frac{b}{2})} \right)^2, \\
 N^{(\text{hyp})} &= \frac{[4 d_{(\text{hyp})}]^b}{2 \Gamma(b) [\zeta(b, \frac{1}{4}) - \zeta(b, \frac{3}{4})]}, & d_{(\text{hyp})} &= \frac{b}{4} \frac{\zeta(1+b, \frac{1}{4}) - \zeta(1+b, \frac{3}{4})}{\zeta(b, \frac{1}{4}) - \zeta(b, \frac{3}{4})}, \\
 \zeta(b, a) &= \sum_{k=0}^{\infty} (k+a)^{-b} \text{ is the generalized Riemann zeta function}
 \end{aligned}$$

TABLE I: The three models used for the shape function forms

The parameters Λ and b are related to the HQET parameters $\bar{\Lambda}$ and μ_π^2 by analytical expressions Eq.46 and Eq.47 in Ref. [5] for exponential and gaussian models, respectively, while for the hyperbolic model the corresponding HQET parameters have to be calculated numerically. The shape function parameters $m_b(\text{SF})$ and $\mu_\pi^2(\text{SF})$ are obtained from the HQET parameters $\bar{\Lambda}$ and μ_π^2 using the relations in Eq.41 of Ref. [5], where the reference scale of 1.5 GeV is used.

Monte Carlo simulated photon energy spectrum

We generate $B \rightarrow X_s \gamma$ MC events according to the prescription in Ref. [5] for each set of the shape function parameter values. The generated events are simulated for the detector performance using the Belle detector simulation program and the $B \rightarrow X_s \gamma$ analysis cuts

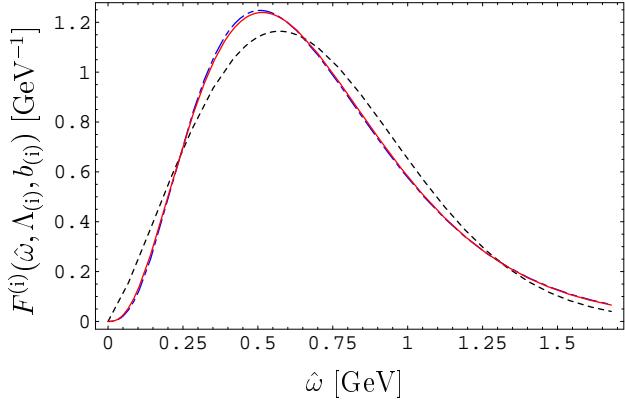


FIG. 1: Three different models with linear onset for small $\hat{\omega}$. We show $F^{(\text{exp})}$ (solid), $F^{(\text{gauss})}$ (dashed), and $F^{(\text{hyp})}$ (dash-dotted), for parameters that correspond to $m_b(\text{SF}) = 4.63 \text{ GeV}/c^2$ and $\mu_\pi^2(\text{SF}) = 0.2 \text{ GeV}^2/c^2$ [5].

are applied to the MC events to obtain the raw photon energy spectrum in the $\Upsilon(4S)$ rest frame [9].

Fitting the spectrum

For a given set of shape function parameters, a χ^2 fit of the MC simulated photon spectrum to the raw data spectrum is performed in the interval $1.8 < E_\gamma^* < 2.8 \text{ GeV}$ [12]. Although in the Ref. [2] the fitting was performed in the interval between 1.5 GeV and 2.8 GeV , the data below 1.8 GeV are not used in the present analysis since the tails of the models we use below 1.8 GeV are not modelled accurately [10].

The normalization parameter is floated in the fit. The raw spectrum is plotted in Figure 2, the errors include both statistical and systematic errors. The latter are dominated by the estimation of the $B\bar{B}$ background and are 100% correlated. Therefore the covariance matrix is constructed as

$$V_{ij} = \sigma_{d_i}^{\text{stat}} \sigma_{d_j}^{\text{stat}} \delta_{ij} + \sigma_{d_i}^{\text{sys}} \sigma_{d_j}^{\text{sys}} \quad (1)$$

where $i, j = 1, 2, \dots, 10$ denote the bin number, and σ_d is the error in the data. Then the χ^2 used in the fitting is given by

$$\chi^2 = \sum_{ij} (d_i - f_i) (V^{-1})_{ij} (d_j - f_j), \quad (2)$$

where $(V^{-1})_{ij}$ denotes the ij^{th} element of the inverted covariance matrix. The χ^2 value of the fit is used to determine a map of χ^2 as a function of the shape function parameters.

The best fit and $\Delta\chi^2$ contour

The best fit parameters are associated to the minimum chi-squared case, χ^2_{\min} . The error “ellipse” is defined as the contour which satisfies $\Delta\chi^2 \equiv \chi^2 - \chi^2_{\min} = 1$. The contours are

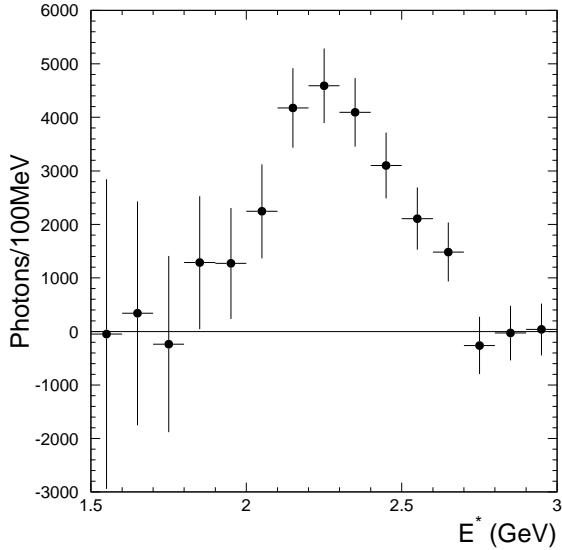


FIG. 2: Raw $B \rightarrow X_s\gamma$ photon energy spectra in the $\Upsilon(4S)$ frame as acquired from data. The errors include both statistical and systematic errors. Raw refers to spectra as measured after the application of Belle $B \rightarrow X_s\gamma$ analysis cuts.

found to be well approximated by the function[11],

$$\Delta\chi^2(m_b(\text{SF}), \mu_\pi^2(\text{SF})) = \left(\frac{\mu_\pi^2(\text{SF}) + a(m_b(\text{SF}))^2 + b}{c} \right)^2 + \left(\frac{(m_b(\text{SF}))^2 + d}{e} \right)^2. \quad (3)$$

The parameters a , b , c , d , and e are determined by fitting the function to the parameter points that lie on the contour.

RESULTS

The best fit parameters are given in Table II. The parameter values are found to be consistent across all three shape function forms. The minimum χ^2 fit for each shape function model is displayed in Figure 3. The fits to the contour with $\Delta\chi^2 = 1$ points are shown in Figure 3 and 4. The imposed shape function form acts to correlate $m_b(\text{SF})$ and $\mu_\pi^2(\text{SF})$.

Shape	χ^2_{min}	$m_b(\text{SF})$ GeV/ c^2	$\mu_\pi^2(\text{SF})$ GeV $^2/c^2$
exponential	4.32	4.52	0.27
gaussian	3.78	4.54	0.25
hyperbolic	4.41	4.52	0.27

TABLE II: The best fit shape function parameter values.

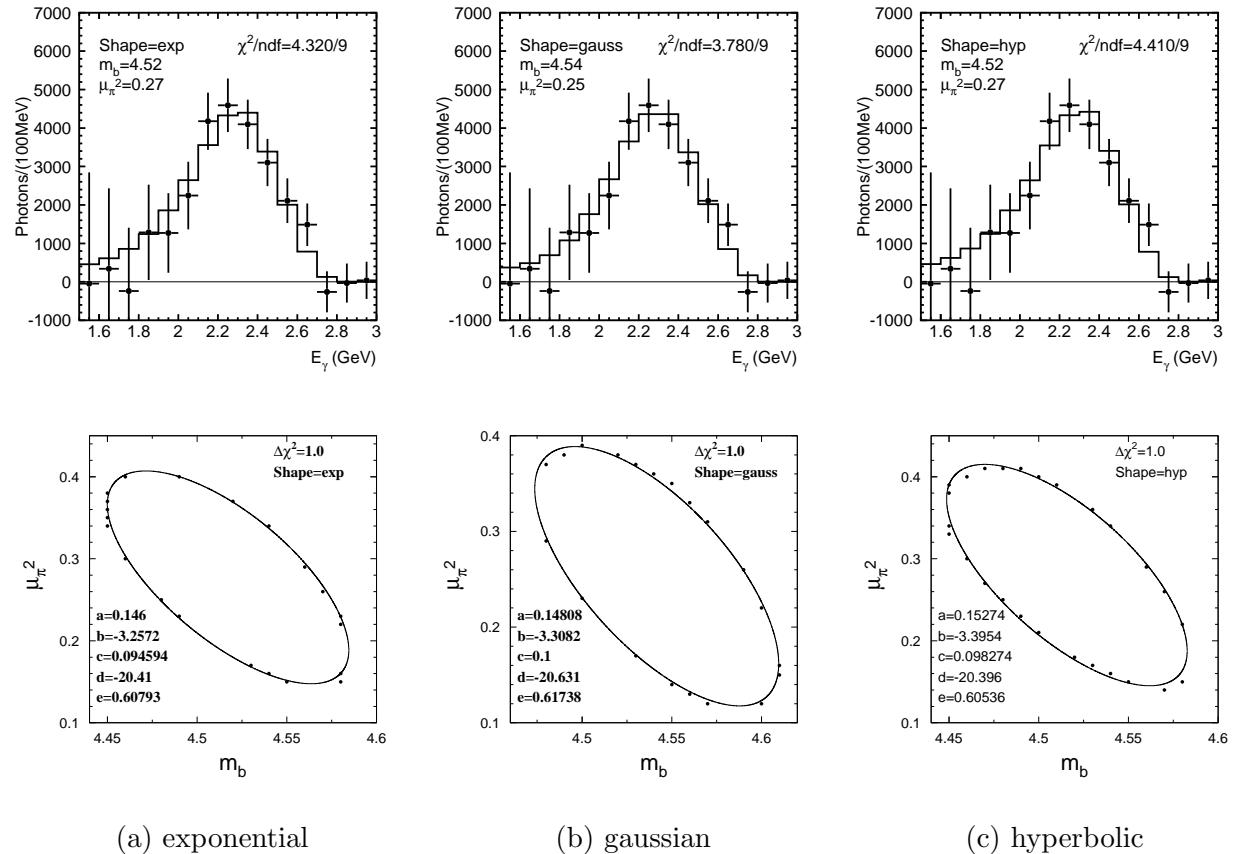


FIG. 3: Top: Minimum χ^2 fits of MC simulated spectra to the raw data for each shape function model. Bottom: The fitted $\Delta\chi^2 = 1$ contours for each shape function model.

SUMMARY

The b -quark leading shape function parameters in the shape function scheme, $m_b(\text{SF})$ and $\mu_\pi^2(\text{SF})$, were determined from fits of Monte Carlo simulated spectra, generated by the prescription in Ref. [5], to the raw [13] Belle measured $B \rightarrow X_s \gamma$ photon energy spectrum. Three models for the form of the leading shape function were used; exponential, gaussian and hyperbolic, while the default model from Ref. [5] was used for the subleading shape function, where the reference scale is chosen to be 1.5 GeV. Best fit parameters are: $(m_b(\text{SF}), \mu_\pi^2(\text{SF}))_{\text{exp}} = (4.52, 0.27)$, $(m_b(\text{SF}), \mu_\pi^2(\text{SF}))_{\text{gauss}} = (4.54, 0.25)$, and $(m_b(\text{SF}), \mu_\pi^2(\text{SF}))_{\text{hyp}} = (4.52, 0.27)$, where $m_b(\text{SF})$ and $\mu_\pi^2(\text{SF})$ are measured in units of GeV/c^2 and GeV^2/c^2 respectively. We also determined the $\Delta\chi^2 = 1$ contours in the $(m_b(\text{SF}), \mu_\pi^2(\text{SF}))$ parameter space for each of the assumed models.

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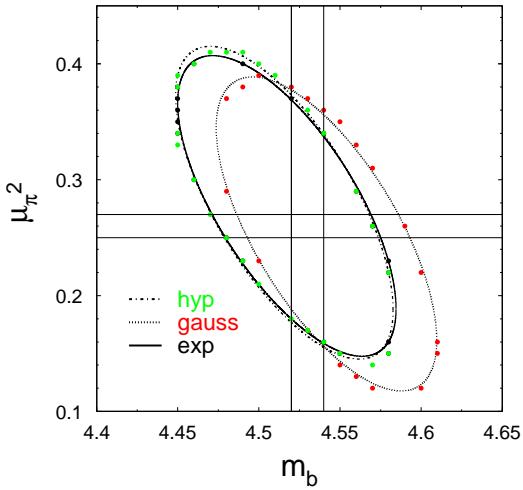


FIG. 4: Comparison of the fitted $\Delta\chi^2 = 1$ contours for all shape function models. The contours for the exponential, gaussian and hyperbolic model are shown by the solid, dotted and dash-dotted curves, respectively. The vertical and horizontal lines mark the central values of the three fits.

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- [11] We thank R. Faccini for suggesting such a function.
- [12] The * denotes the center of mass frame or equivalently the $\Upsilon(4S)$ rest frame.
- [13] Raw refers to the spectrum as measured after the application of analysis cuts.